

# Scaling in ISiS Au multifragmentation data

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Evidence of thermal scaling was found in the ISiS Au multifragmentation data [1]. Thermal scaling is the feature that the fragment yields depend on temperature  $T$  via a Boltzmann factor:  $n \propto \exp(-B/T)$ ;  $B$  is the fragment emission barrier. A plot of  $\ln n$  vs.  $1/T$  (Arrhenius plot) is linear if  $n$  is controlled by a Boltzmann factor. This is observed in a plot of the log of the fragment yield distributions  $n(A)$  as a function of  $1/\sqrt{E^*}$  ( $\sqrt{E^*} \propto T$  in a Fermi gas) in Fig. 1a.

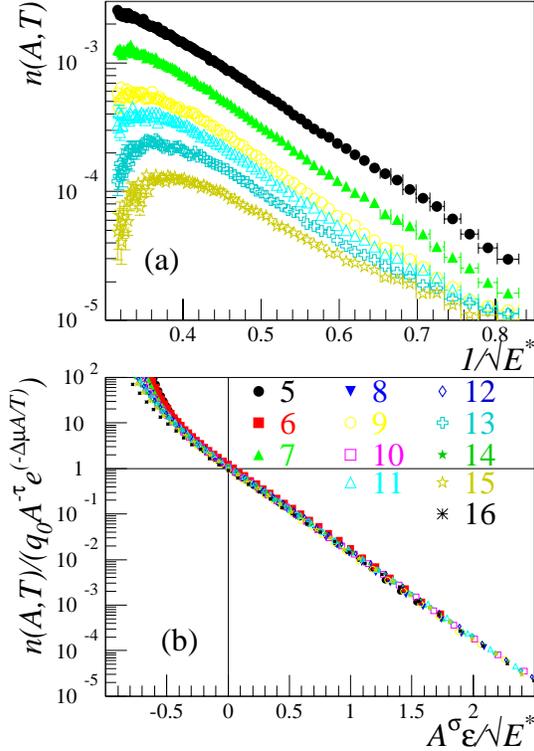


Figure 1: (a) Arrhenius plots. (b) Fisher's plot.

Similar behavior was observed in the EOS Au multifragmentation data [2] and the Boltzmann

factors (slopes of Fig. 1a) were interpreted via Fisher's droplet model [3]:

$$n(A, T) = q_0 A^{-\tau} \exp(A\Delta\mu/T - c_0 \epsilon A^\sigma / T), \quad (1)$$

( $\epsilon = (T_c - T)/T_c$ ) where it was assumed that the system was at coexistence ( $\Delta\mu \approx 0$ ). This lead to a relation between  $B$  and fragment mass  $A$ :  $B = c_0 A^\sigma$ ;  $c_0$  is the surface energy coefficient. Within the modest statistics of the EOS data set ( $\sim 27K$  events) the  $\Delta\mu \approx 0$  assumption was difficult to test. Using the high statistics of the ISiS data set ( $\sim 7.4M$  events) it is now possible.

The ISiS yields for  $E^* \leq E_c^*$  and  $5 \leq Z \leq 16$  were fit to Eq. (1) using  $A$  determined by multiplying the measured  $Z$  by the  $A$ -to- $Z$  ratio of the fragmenting system and  $T \propto \sqrt{E^*}$ ;  $\tau$ ,  $\sigma$ ,  $c_0$  and  $E_c^*$  were fit paramters and  $\Delta\mu$  was parametrized by a polynomial in  $E^*$ . The results are:  $\tau = 2.03 \pm 0.01$ ,  $\sigma = 0.65 \pm 0.04$ ,  $E_c^* = 3.70 \pm 0.05$  MeV/nucleon,  $c_0 \sqrt{a} \approx 4.25 \pm 0.05$  and  $\Delta\mu$  decreased from  $\sim 1$  MeV at low  $E^*$  to 0 at  $E_c^*$ .

Fisher's model predicts that a graph of the scaled fragment yields will collapse the data from a wide range of mass and temperature collapse onto a single curve, thereby providing the best evidence yet for a liquid-vapor phase transition in nuclei (Fig. 1b).

## References

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- [2] J. B. Elliott *et al.*, Phys. Rev. Lett. **85**, 1194 (2000).
- [3] M. E. Fisher, Physics **3**, 255 (1967).